

## CORRECTIONS

**Y. Wang, S. Ge, M. Rafailovich,\* J. Sokolov, Y. Zou, H. Ade, J. Lüning, A. Lustiger, and G. Marom:** Crystallization in the Thin and Ultrathin Films of Poly(ethylene–vinyl acetate) and Linear Low-Density Polyethylene. . Volume 37, Number 9, May 4, 2004, pp 3319–3327.

Page 3319. One of the authors' names was incorrectly printed as G. Maron. The correct spelling is G. Marom.

Page 3323. Discussion about the TEM image in Figure 8. As for the decorated fibrous crystal networks in between the big crystal lamellae, although these appear to look like the classic “shish-kebab” crystal, the lack of significant stress during the formation of the structure made it unlikely that these are real shish-kebabs, as tentatively suggested in the paper. It is more probable that they are immature lamellae formed by the chains with lower molecular weight which were rejected from the larger crystalline lamellae during crystallization.

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**Félix S. Csajka and Christian Seidel\*:** Strongly Charged Polyelectrolyte Brushes: A Molecular Dynamics Study. . Volume 33, Number 7, April 4, 2000, pp 2728–2739.

Due to a numerical error, unfortunately the simulations were effectively carried out at a Bjerrum length different from that given in the paper. In fact collapsed brushes were not simulated at  $\lambda_B = 0.7\sigma$  but instead at  $\lambda_B \approx 14\sigma$ . In addition to that, the numerical error causes a weak dependence of the real Bjerrum length on grafting density. In the collapsed state, however, this has no remarkable effect on the stretching of the chains. Note that the collapsed brush regime is not so exotic as it might appear due to the large Bjerrum length reported here. From theory, we know that the existence of this particular brush regime crucially depends on the relation between second virial and Bjerrum length.<sup>1</sup> While many polyelectrolytes have a hydrophobic backbone, in our model, we assume simple good solvent behavior with a rather large second virial. Therefore the collapsed regime is shifted to large  $\lambda_B$ .

## References and Notes

- (1) Csajka, F. S.; Netz, R. R.; Seidel, C.; Joanny, J.-F. *Eur. Phys. J. E.* **2001**, *4*, 505.

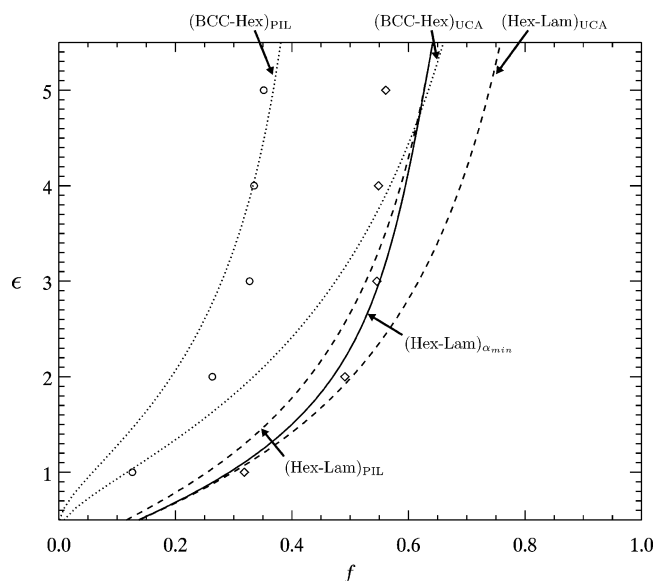
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**Gregory M. Grason and Randall D. Kamien\*:** Interfaces in Diblocks: A Study of Miktoarm Star Copolymers. . Volume 37, No. 19, September 21, 2004, pp 7371–7380.

The labels of the curves in Figure 8 are incorrect. Below is the correct labeling and caption for Figure 8.



**Figure 8.** Dashed line shows the predicted Lam–Hex phase boundaries computed using the PIL, upper-bound (leftmost boundary), and UCA, lower-bound (rightmost boundary), approximations for the SST Hex phase. The solid line is the Lam–Hex boundary computed using the SST results of section III. The dotted lines are the Hex–BCC phase boundaries by comparing the PIL, upper-bound (leftmost), and UCA, lower-bound (rightmost), SST results for those morphologies. For comparison the  $\chi N \approx 100$  SCFT results for the Lam–Hex boundary are shown as diamonds,  $\diamond$ , and the BCC–Hex boundary are shown as open circles,  $\circ$ .

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